



دانشگاه صنعتی اصفهان
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Elastic–Plastic Fracture Mechanics



- Linear elastic fracture mechanics (LEFM) is valid only as long as nonlinear material deformation is confined to a small region surrounding the crack tip. In many materials, it is virtually impossible to characterize the fracture behavior with LEFM, and an alternative fracture mechanics model is required.
- Elastic–plastic fracture mechanics applies to materials that exhibit time-independent, nonlinear behavior (i.e., plastic deformation). Two elastic–plastic parameters are introduced: the crack tip opening displacement (*CTOD*) and the *J integral*. Both parameters describe crack tip conditions in elastic–plastic materials, and each can be used as a fracture criterion. Critical values of CTOD or J give nearly size-independent measures of fracture toughness, even for relatively large amounts of crack tip plasticity.

CTOD: Crack Tip Opening Displacement

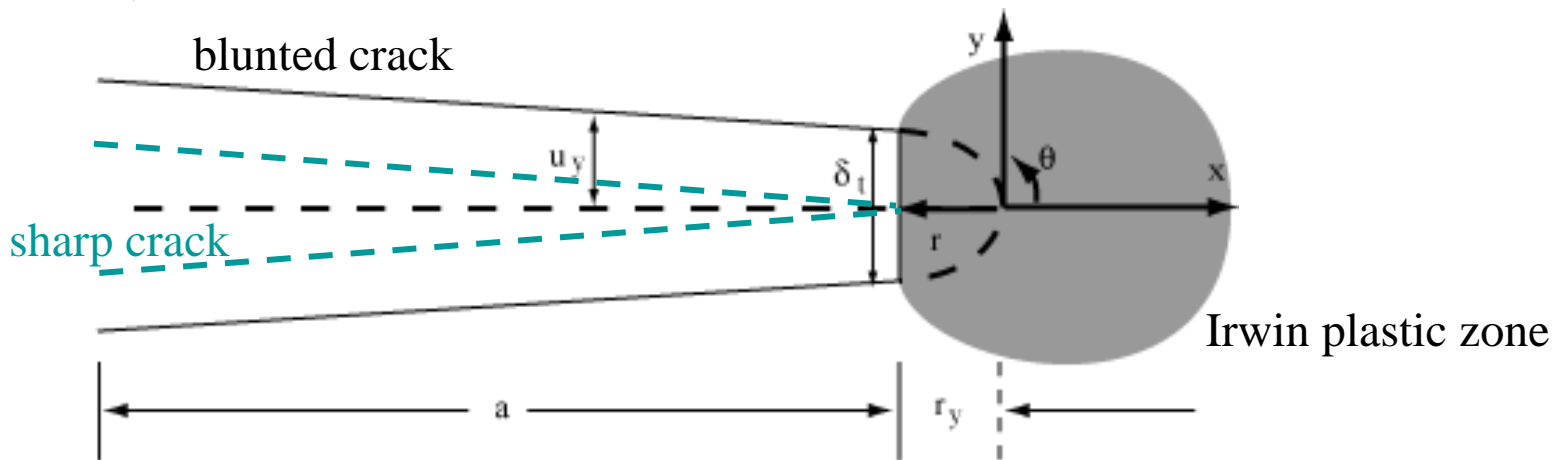
Wells's experimental work: attempt to measure K_{Ic} for structural steels



But

Initial sharp crack has *blunted* prior to fracture

Non-negligible plastic deformation



➔ LEFM *inaccurate* : materials too tough !!!

Instead, Wells proposed δ_t (CTOD) as a measure of fracture toughness.

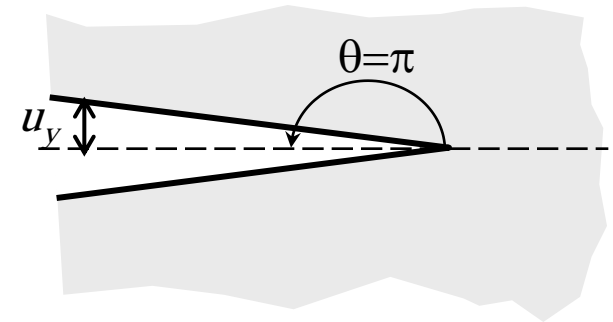
Estimation of δ_t using Irwin model : Crack length: $a + r_y$

By definition, $\delta_t = 2u_y$ at $r = r_y$ where u_y is the crack opening

Crack opening: (u_y)

$$u_y = \frac{K_I}{2\mu} \sqrt{\frac{r}{2\pi}} \sin \frac{\theta}{2} \left[\kappa + 1 - 2 \cos^2 \frac{\theta}{2} \right] \Big|_{\theta=\pi}$$

$$= \frac{K_I}{2\mu} (\kappa + 1) \sqrt{\frac{r}{2\pi}} \quad (\text{see Table 2.2})$$



We have $\mu = \frac{E}{2(1+\nu)}$ and for plane stress, $\kappa = \frac{3-\nu}{1+\nu}$

$$\Rightarrow \delta_t = 2 \frac{4K_I}{E} \sqrt{\frac{r_y}{2\pi}}$$

From Irwin model, the radius of the plastic zone is $r_y = \frac{1}{2\pi} \left(\frac{K_I}{\sigma_{YS}} \right)^2$

$$\delta_t = \frac{4}{\pi} \frac{K_I^2}{\sigma_{YS} E} \quad \text{and also,} \quad \delta_t = \frac{4}{\pi} \frac{G}{\sigma_{YS}}$$

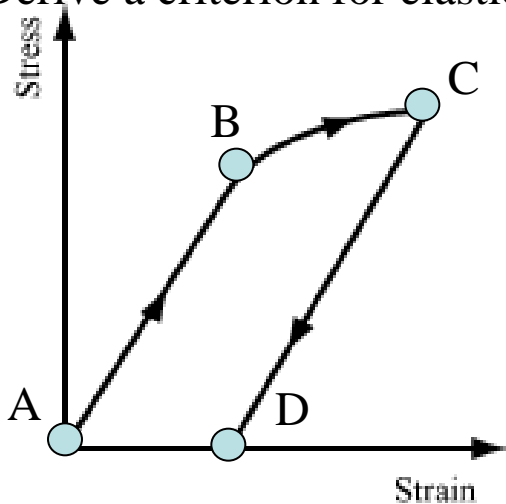
CTOD related uniquely to K_I and G .

➔ CTOD appropriate characterizing crack-tip-parameter when LEFM no longer valid.

Can be proved by a unique relationship between CTOD and the J integral.

The J contour integral as yield criterion

- More general criterion than K (valid for LEFM)
- Derive a criterion for elastic-plastic materials, with typical stress-strain behavior:



A→B : linear

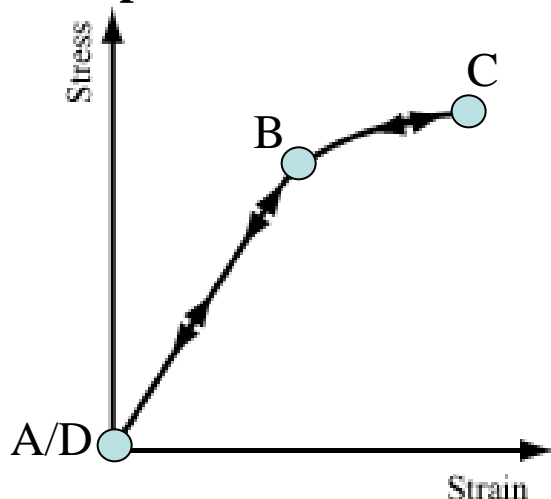
B→C : non-linear curve

C→D : non-linear, same slope as A-B

non-reversibility: $A-B-C \neq C-D-A$

- Material behavior is **strain history dependent** !
Non unique solutions for stresses

- Simplification:** non-linear elastic behavior



reversibility: $A-B-C = C-D-A$

- Correct **only** for a **monotonic** loading

= Deformation theory of Plasticity



The J contour integral as yield criterion

Definition of the J-integral

Rice defined a *path-independent* contour *integral J* for the analysis of cracked bodies showed that its value = *energy release rate in a nonlinear elastic material*

J generalizes *G* to nonlinear materials :

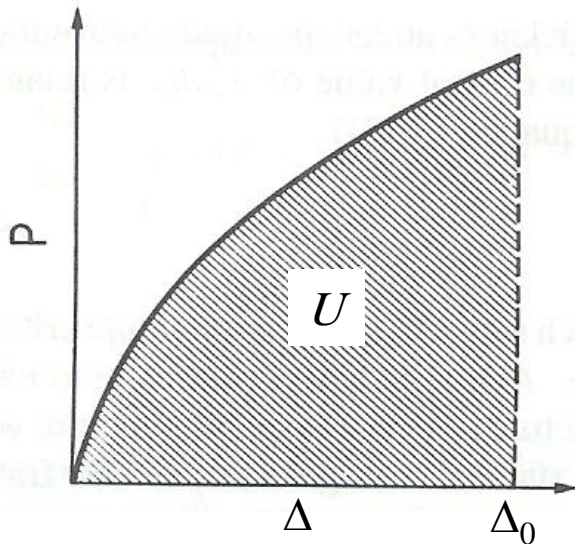
→ nonlinear elastic energy release rate

As *G* can be used as a fracture criterion *J_c*

reduces to *G_c* in the case of linear fracture

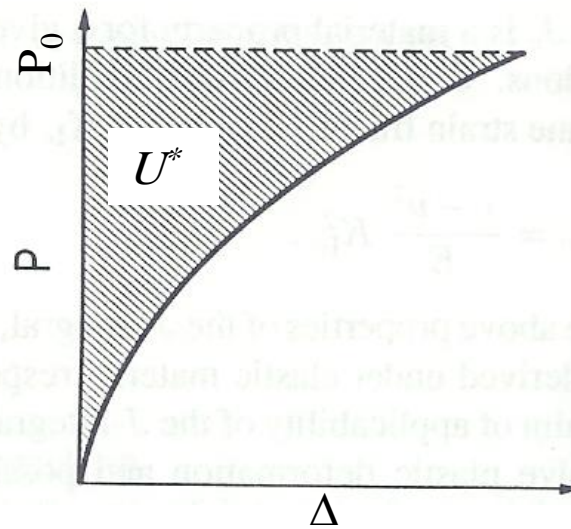
Definition of the J-integral

- Historically,
 - Rice defined a *path-independent* contour *integral* J for the analysis of crack
 - showed that its value = *energy release rate* in a *nonlinear* elastic body with a crack
- J generalizes the concept of G to non-linear materials
 - For linear materials $J = G$
 - Load-displacement diagram: potential energy Π



Fixed-grips conditions:

$$\Pi = U = \int_0^{\Delta_0} P(\Delta) d\Delta$$



Dead-load conditions

$$-\Pi = U^* = \int_0^{P_0} \Delta(P) dP$$

U : Elastic strain energy

\neq (in general)

U^* : Complementary energy

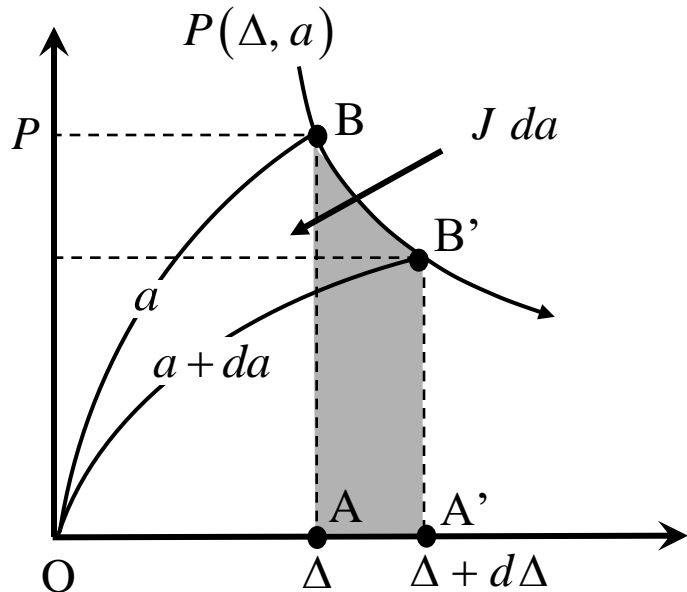
The J contour integral as yield criterion

- Definition of J using the potential energy Π :

$$J = -\frac{d\Pi}{dA}$$

$A = a B$: for a cracked plate with through crack

- Geometrical interpretation:



OB and OB' :

loading/unloading for the given body with crack lengths a and $a+da$

$P(\Delta, a)$:

Possible relationship between the load P and the displacement Δ while the crack is moving.

We have $J dA = Pd\Delta - dU$

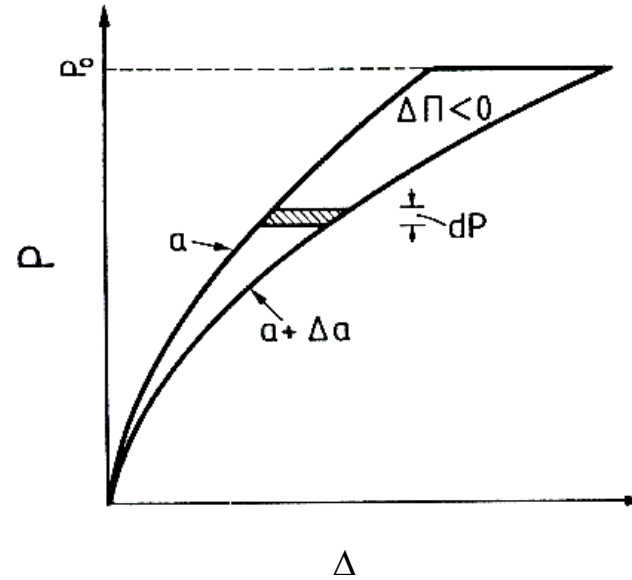
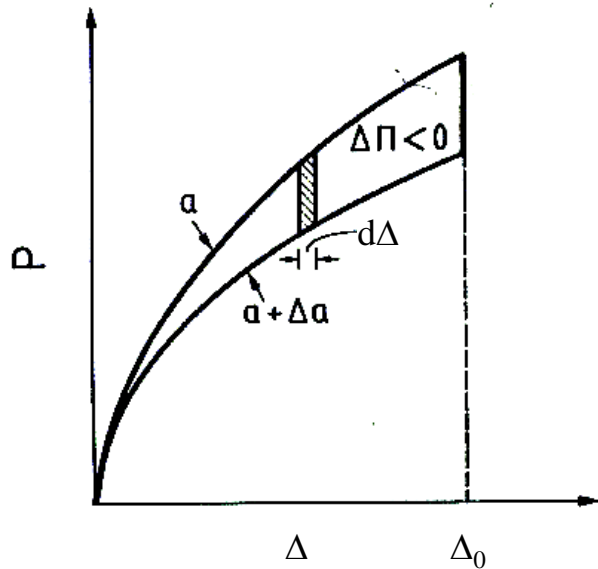
dU is the difference between the areas under OB' and OB : $OA'B' - OAB$

$Pd\Delta$ appears as the area $AA'B'B$

Thus, $J dA = J B da = AA'B'B + OAB - OA'B' = OBB'$

The J contour integral as yield criterion

- In particular ,



At constant displacement:

$$J = -\frac{1}{B} \left(\frac{\partial U}{\partial a} \right)_{\Delta} = -\frac{1}{B} \int_0^{\Delta_0} \left(\frac{\partial P}{\partial a} \right)_{\Delta} d\Delta$$

At constant force (dual form):

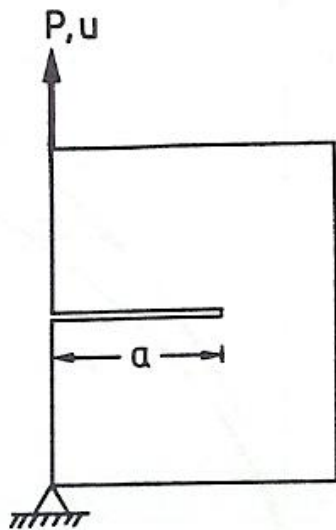
$$J = \frac{1}{B} \left(\frac{\partial U^*}{\partial a} \right)_P = \frac{1}{B} \int_0^{P_0} \left(\frac{\partial \Delta}{\partial a} \right)_P dP$$

Useful expressions for the experimental determination of J

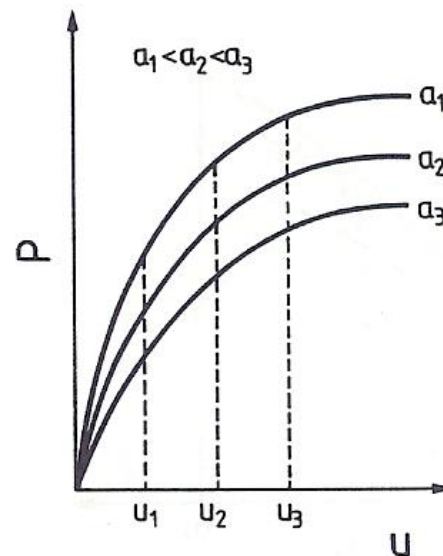
- Experimental determination of the J-integral :
 - Multiple-specimen method (Begley and Landes (1972)) :

Procedure

(1) Consider cracked specimens with different crack lengths a_i



(1)

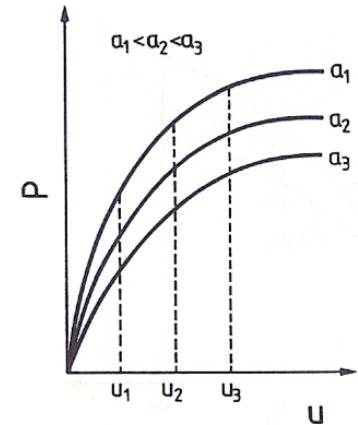


(2)

The J contour integral as yield criterion

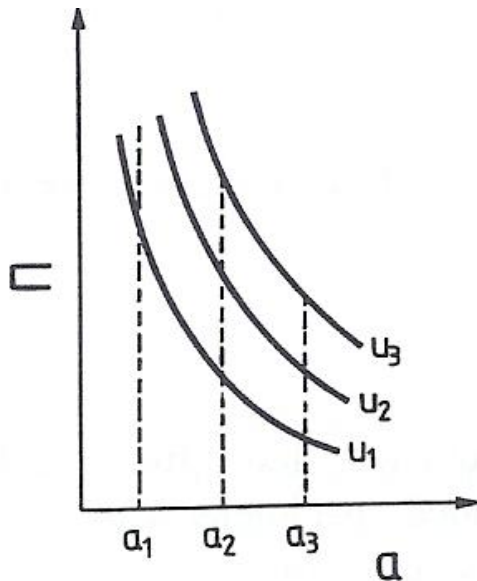
(3) Calculation of the potential energy Π for given values of displacement u

= area under the load-displacement curve

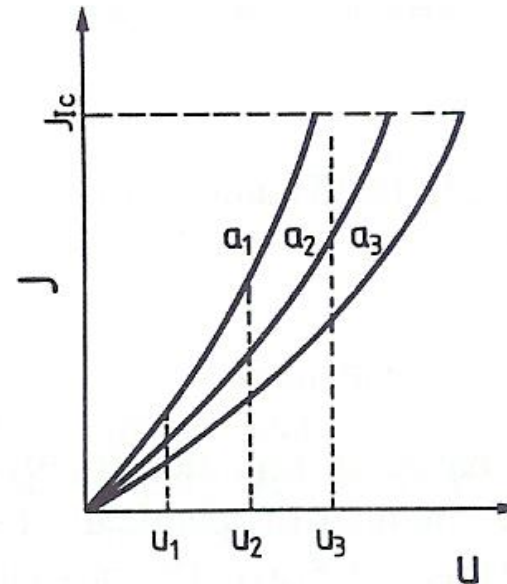


(4) Negative slopes of the $P - a$ curves determined and plotted versus displacement for different crack lengths :

Critical value J_{Ic} of J at the onset of crack extension (material constant)



(3)



(4)

The J contour integral as yield criterion

- J as a path-independent line integral

$$J = \int_{\Gamma} \left(w \, dy - T_i \frac{\partial u_i}{\partial x} \, ds \right) \quad \text{with} \quad w(\varepsilon_{mn}) = \int_0^{\varepsilon_{mn}} \sigma_{ij} \, d\varepsilon_{ij} \quad \text{strain energy density}$$

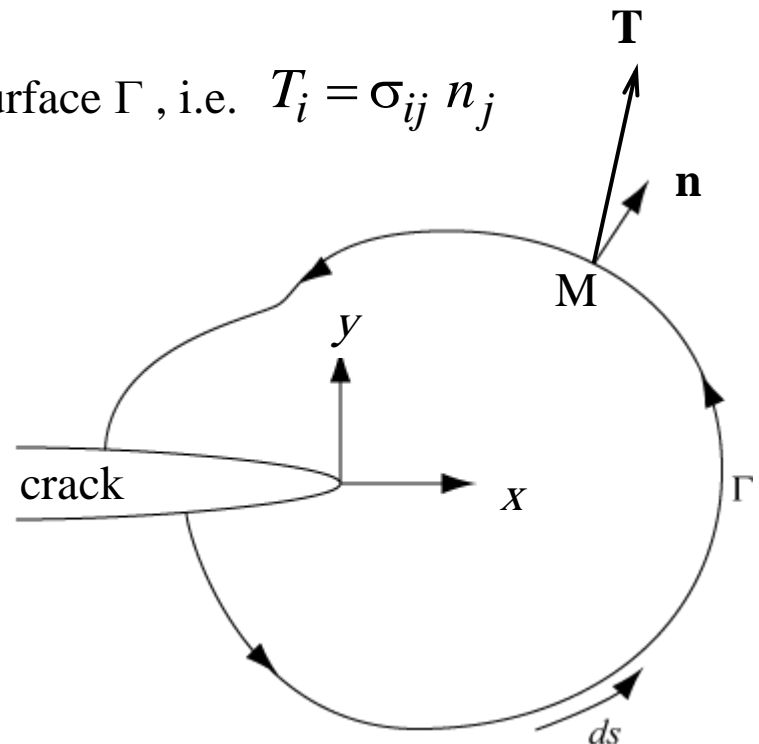
$$= \int_{\Gamma} \left(w \, dy - \mathbf{T} \cdot \frac{\partial \mathbf{u}}{\partial x} \, ds \right)$$

\mathbf{T} : traction vector at a point M on the bounding surface Γ , i.e. $T_i = \sigma_{ij} n_j$

\mathbf{u} : displacement vector at the same point M.

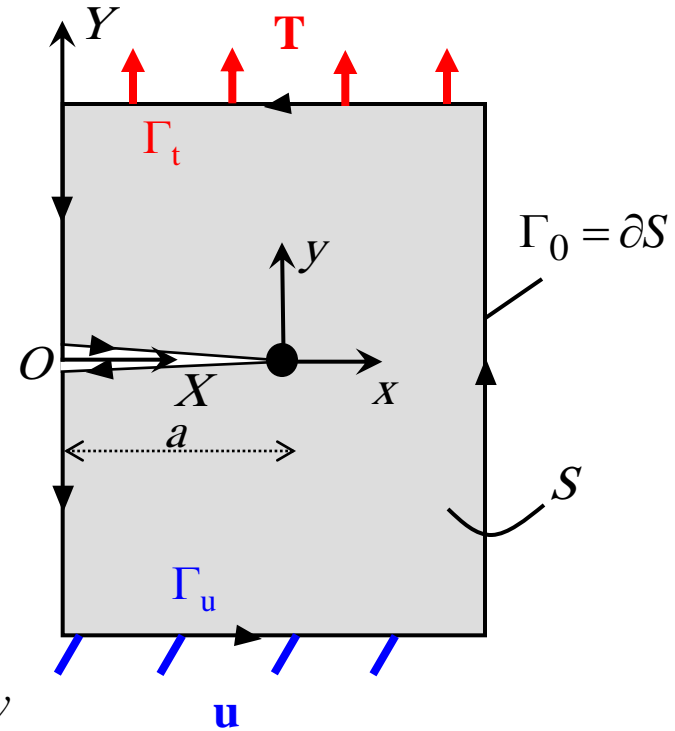
\mathbf{n} : unit *outward* normal.

The contour Γ is followed in the *counter-clockwise* direction.



Equivalence of the two definitions

- 2D solid of unit thickness of area S ,
with a linear crack of length a along OX (fixed)
- Crack faces are traction-free.
- Total contour of the solid Γ_0 *including* the crack tip:
Imposed **tractions** on the part of the contour Γ_t
Displacements applied on Γ_u



Proof : Recall for the potential energy (per unit thickness),

$$\Pi(a) = \iint_S w dS - \int_{\Gamma_t} T_i u_i ds \quad T_i = \sigma_{ij} n_j \quad \sigma_{ij} = \frac{\partial w}{\partial \varepsilon_{ij}}$$

The tractions and displacements imposed on Γ_t and Γ_u are independent of a

$$\begin{aligned} \frac{dT_i}{da} &= 0, \quad \text{on } \Gamma_t \\ \frac{du_i}{da} &= 0 \quad \text{on } \Gamma_u \end{aligned} \quad \Rightarrow \quad \frac{d\Pi}{da} = \iint_S \frac{dw}{da} dS - \int_{\Gamma_0} T_i \frac{du_i}{da} ds$$



The J contour integral as yield criterion

Considering the moving coordinate system x, y (attached to the crack tip), $x = X - a$

$\frac{d}{da}$: total derivative/crack length

$$\frac{d}{da} = \left(\frac{\partial}{\partial a} \right)_x + \left(\frac{\partial x}{\partial a} \right)_X \left(\frac{\partial}{\partial x} \right)_a = \frac{\partial}{\partial a} - \frac{\partial}{\partial x}$$

Thus,

$$\frac{d\Pi}{da} = \iint_S \left(\frac{\partial w}{\partial a} - \frac{\partial w}{\partial x} \right) dS - \int_{\Gamma_0} T_i \left(\frac{\partial u_i}{\partial a} - \frac{\partial u_i}{\partial x} \right) ds$$

However,

$$\begin{aligned} \frac{\partial w}{\partial a} &= \frac{\partial w}{\partial \varepsilon_{ij}} \frac{\partial \varepsilon_{ij}}{\partial a} = \sigma_{ij} \frac{\partial \varepsilon_{ij}}{\partial a} = \sigma_{ij} \frac{\partial}{\partial a} \left[\frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \right] && \text{since } \sigma_{ij} = \sigma_{ji} \\ &= \sigma_{ij} \frac{\partial}{\partial a} \frac{\partial u_i}{\partial x_j} = \sigma_{ij} \frac{\partial}{\partial x_j} \left(\frac{\partial u_i}{\partial a} \right) \end{aligned}$$



The J contour integral as yield criterion

Thus,

$$\iint_S \frac{\partial w}{\partial a} dS = \iint_S \sigma_{ij} \frac{\partial}{\partial x_j} \left(\frac{\partial u_i}{\partial a} \right) dS$$

We have,

$$\iint_S \sigma_{ij} \frac{\partial}{\partial x_j} \left(\frac{\partial u_i}{\partial a} \right) dS = \int_{\Gamma_0} \sigma_{ij} \frac{\partial u_i}{\partial a} n_j ds = \int_{\Gamma_0} T_i \frac{\partial u_i}{\partial a} ds$$

The derivative of J reduces to,

$$\begin{aligned} \frac{d\Pi}{da} &= \iint_S \left(\frac{\partial w}{\partial a} - \frac{\partial w}{\partial x} \right) dS - \int_{\Gamma_0} T_i \left(\frac{\partial u_i}{\partial a} - \frac{\partial u_i}{\partial x} \right) ds \\ &= -\iint_S \left(\frac{\partial w}{\partial x} \right) dS + \int_{\Gamma_0} T_i \left(\frac{\partial u_i}{\partial a} \right) ds - \int_{\Gamma_0} T_i \left(\frac{\partial u_i}{\partial a} - \frac{\partial u_i}{\partial x} \right) ds \\ &= -\left(\iint_S \left(\frac{\partial w}{\partial x} \right) dS - \int_{\Gamma_0} T_i \left(\frac{\partial u_i}{\partial x} \right) ds \right) \end{aligned}$$



The J contour integral as yield criterion

Using the Green Theorem, i.e. $\oint_{\Gamma} P(x, y) dx + Q(x, y) dy = \iint_A \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dx dy$

$$\begin{aligned} -\frac{d\Pi}{da} &= \iint_S \left(\frac{\partial w}{\partial x} \right) dS - \int_{\Gamma_0} T_i \left(\frac{\partial u_i}{\partial x} \right) ds \\ &= \int_{\Gamma_0} \left(w dy - T_i \left(\frac{\partial u_i}{\partial x} \right) ds \right) \end{aligned}$$

➡ J derives from a potential